

## Project Aims and Challenges

This project focuses in the derivation of **Kinetic-Induced Moment Systems (KIMS)** based in the relationship between kinetic theory and non-linear hyperbolic conservation laws and their potential applications in the description of small-scale geophysical flows. Using simple base cases as the 1-D Burgers' equation and the 1-D shallow water equations, our aim on the one hand is to use the resulting PDEs as a monitoring function to detect particular flow structure (like shocks and rarefaction waves) into the construction of adaptive numerical methods, and on the other hand as a basis to derive novel parametrizations for subgrid closures.

## Main Definitions

- **Boltzmann-like equation with a BGK collision term:** describes the statistical distribution of the density of particles  $f(t, x, \xi)$ , where  $f_0(Q, \xi)$  is considered its *equilibrium function* and  $\xi$  the microscopic velocity,

$$\partial_t f(t, x, \xi) + \xi \partial_x f(t, x, \xi) = \frac{1}{\varepsilon} [f_0(Q, \xi) - f(t, x, \xi)]$$

with  $0 < \varepsilon \ll 1$ , the mean free path.

- **Moments ( $W_k$ ):** are weighted averages of  $f(t, x, \xi)$ ,

$$W_k = \int_{\mathbb{R}} \xi^k f(t, x, \xi) d\xi$$

with  $k \in \mathbb{Z}_{\geq 0}$  the order of the moment.

The **equilibrium moments**  $W_k|_E$  are also weighted averages but of  $f_0(Q, \xi)$ ,

$$W_k|_E = \int_{\mathbb{R}} \xi^k f_0(Q, \xi) d\xi$$

Consider a 1-D conservation laws system

$$\partial_t Q + \partial_x F(Q) = 0$$

$Q(n)$  the vector of unknowns with  $n$  elements. For  $0 < k \leq n-1$  holds that  $f(t, x, \xi) = f_0(Q, \xi)$  and

$$W_{k-1} = W_{k-1}|_E = Q(k), \quad 0 < k \leq n-1$$

and the remaining moments  $W_k$  for  $k \geq n$ , will yield the new unknowns.

## What is a KIMS?

It is an infinite moment system based on an artificial Boltzmann-like transport equation using the **connection between kinetic theory and conservation laws** together with an asymptotic expansion of the corresponding moments.

## How to derive it?

1. Multiplication on both sides of the Boltzmann-like transport equation by the weights  $(1, \xi, \dots, \xi^{n-1}, \xi^k)$  and subsequent integration over the microscopic velocity  $\xi$  yield an infinite PDE moment system.

2. Express the infinite system in terms of the  $p$ -order non-equilibrium moments, starting at  $p = 1$

$$W_k^{(p)} = W_k^{(p-1)} - W_k^{(p-1)}|_E, \quad p \geq 1$$

for  $0 \geq k \geq n+p-2$  holds that  $W_k^{(p)} = 0$ .

3. Use asymptotic expansion in terms of the small parameter  $\varepsilon$  to show that at each order a scale-induced closure is possible, resulting in a closed moment system.

$$W_k^{(p)} = \varepsilon^p W_{k,p}^{(p)} + \varepsilon^{p+1} W_{k,p+1}^{(1)} + \dots, \quad k \geq n+p-1$$

## Base Cases

1. 1-D Inviscid Burgers' equation:

$$\partial_t u + \partial_x \left( \frac{1}{2} u^2 \right) = 0$$

evolution of the horizontal velocity  $u(t, x)$  at a time  $t \geq 0$  and at a point  $x \in \mathbb{R}$ .

2. 1-D Shallow Water equations:

$$\begin{aligned} \partial_t h + \partial_x(hu) &= 0 \\ \partial_t(hu) + \partial_x(hu^2 + \frac{g}{2}h^2) &= 0 \end{aligned}$$

evolution of the height of water  $h(t, x)$  and its horizontal velocity  $u(t, x)$ , with  $t \geq 0$ ,  $x \in \mathbb{R}$  and  $g$  the gravitational acceleration.

The respective KIMS at third order ( $p = 3$ ) reads,

- 1-D Inviscid Burgers' equation:

$$\begin{aligned} \partial_t u + u \partial_x u + \varepsilon \partial_x W &= 0 \\ \partial_t W + \frac{1}{3\varepsilon} \partial_x u + u \partial_x W &= \frac{4\varepsilon}{15} \partial_{xx} W - \frac{1}{\varepsilon} W \end{aligned}$$

where  $W = \frac{1}{\varepsilon} W_1^{(1)}$

- 1-D Shallow Water equations:

$$\begin{aligned} \partial_t h + \partial_x(hu) &= 0 \\ \partial_t(hu) + \partial_x(hu^2 + \frac{g}{2}h^2) + \varepsilon \partial_x W &= 0 \\ \partial_t W + \left( \frac{g}{2\varepsilon} h^2 + 3W \right) \partial_x u + u \partial_x W &= -\frac{3}{2} \varepsilon g \partial_x \left[ \frac{g}{\varepsilon} h^2 \partial_x h + 2W \partial_x h + \right. \\ &\quad \left. h \partial_x W + gh^3 \partial_{xx} u \right] - \frac{1}{\varepsilon} W \end{aligned}$$

where  $W = \frac{1}{\varepsilon} W_2^{(1)}$

## Monitoring Functions ( $\varepsilon \rightarrow 0$ )

In the limit  $\varepsilon \rightarrow 0$ , the previous third-order systems read

- 1-D Inviscid Burgers' equation:

$$\begin{aligned} \partial_t u + u \partial_x u &= 0 \\ W &= -\frac{1}{3} \partial_x u \end{aligned}$$

- 1-D Shallow Water equations:

$$\begin{aligned} \partial_t h + \partial_x(hu) &= 0 \\ \partial_t(hu) + \partial_x(hu^2 + \frac{g}{2}h^2) &= 0 \\ W &= -\frac{g}{2} h^2 \partial_x u \end{aligned}$$

Can be proved that the previous systems yield the correct shock propagation and therefore as  $\varepsilon \rightarrow 0$ : if  $\partial_x u = 0$  then  $W = 0$ , if  $\partial_x u \rightarrow -\infty$  then  $W \rightarrow \infty$  (shock wave) and if  $\partial_x u \rightarrow \infty$  then  $W \rightarrow -\infty$  (rarefaction wave). Consequently,  $W(x, t)$  tends to  $\delta$ -function located at the points of the discontinuities.

## Numerical Experiments

Using developing shock initial conditions in both cases, we proof numerically that for a sufficiently small epsilon ( $\varepsilon = 0.01$ ),  $W(x, t)$  behaves as expected.

- 1-D Inviscid Burgers' equation:

$$u(x, 0) = \tan^{-1}(-x) + 4 \quad \forall x$$

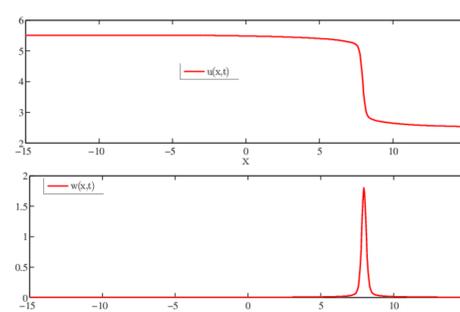


Figure 1.  $u(x, t)$  and  $w(x, t)$  at  $t = 2$

- 1-D Shallow Water equations:

(Dam break)

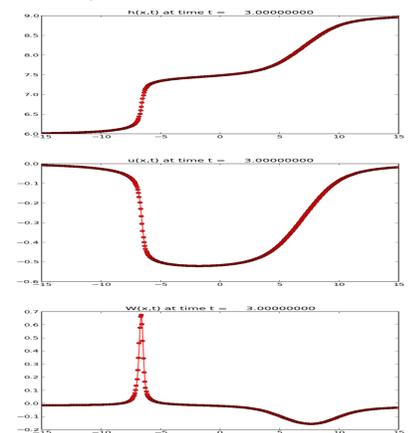


Figure 2.  $h(x, t)$ ,  $u(x, t)$  and  $w(x, t)$  at  $t = 3$

## Present and Future Work

- The current stage of our research is focused on the applicability of the monitoring function as a refinement parameter in the construction of novel grid-adaptive simulation tools. Again, we use the previous base cases in the development of numerical experiments and compare its performance with traditional grid-adaptive simulations.
- The next step consist in the derivation of novel parametrizations for subgrid closures. We will use spectral analysis in order to compare in a coherent way the original flow equations and its corresponding moment system, together with the corresponding subgrid closures.

## Selected References

- [1] G. Diana and J. Struckmeier: *Kinetic-Induced Moment Systems for the Saint-Venant Equations*, TASK QUARTERLY, 2013, Vol. 17 No. 1-2, pp. 6390.
- [2] G. Diana, H. Struchtrup and J. Struckmeier: *A Kinetic-Induced Moment System for Burgers Equation*, 2014 (in preparation).
- [3] B. Perthame: *Kinetic Formulation of Conservation Laws*, OXFORD University Press, NY, 2002.
- [4] B. Perthame and C. Simeoni: "A Kinetic Scheme for the Saint-Venant System with a source term", CALCOLO, 2001, No. 38, pp. 201-231.
- [5] H. Struchtrup: *Macroscopic Transport Equations for Rarefied Gas Flows*, Springer-Verlag, Berlin, 2005.

## Affiliations

<sup>1</sup> Department of Mathematics, University of Hamburg, Germany